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# **OPEN** A hole inversion layer at the BiVO<sub>4</sub>/Bi<sub>4</sub>V<sub>2</sub>O<sub>11</sub> interface produces a high tunable photovoltage for water splitting

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The conversion of solar energy into hydrogen fuel by splitting water into photoelectrochemical cells (PEC) is an appealing strategy to store energy and minimize the extensive use of fossil fuels. The key requirement for efficient water splitting is producing a large band bending (photovoltage) at the semiconductor to improve the separation of the photogenerated charge carriers. Therefore, an attractive method consists in creating internal electrical fields inside the PEC to render more favorable band bending for water splitting. Coupling ferroelectric materials exhibiting spontaneous polarization with visible light photoactive semiconductors can be a likely approach to getting higher photovoltage outputs. The spontaneous electric polarization tends to promote the desirable separation of photogenerated electron-hole pairs and can produce photovoltages higher than that obtained from a conventional p-n heterojunction. Herein, we demonstrate that a hole inversion layer induced by a ferroelectric Bi<sub>4</sub>V<sub>2</sub>O<sub>11</sub> perovskite at the n-type BiVO<sub>4</sub> interface creates a virtual p-n junction with high photovoltage, which is suitable for water splitting. The photovoltage output can be boosted by changing the polarization by doping the ferroelectric material with tungsten in order to produce the relatively large photovoltage of 1.39V, decreasing the surface recombination and enhancing the photocurrent as much as 180%.

The strategy based on storing solar energy in the form of chemical energy is the envisaged development of technology in an attempt to prevent or minimize harmful effects on the natural environmental, which has been recognizably accelerated by the large current use of fuels derived from fossil sources. Several recently reported studies have been focusing on the development of water-splitting photoelectrochemical (PEC) cells, which consists of devices assembled with semiconductor photoelectrodes that operate by using the sunlight radiation for producing H<sub>2</sub> and O<sub>2</sub> through the molecular splitting of water<sup>1,2</sup>. An input energy of 1.23 V is thermodynamically required for such a water splitting reaction and photoelectrodes with bandgap energies below 2.5 eV are needed for an efficient use of sunlight radiation3. To meet the best working conditions, the band edges potentials at the semiconductor surfaces must straddle the H2 and O2 redox potentials, the electrodes should be chemically stable regarding the photocorrosion, and the charge transfer from the surface of the semiconductor to the solution must be efficient enough to minimize the energy losses due to kinetics overpotentials<sup>3,4</sup>. To date, no known single material has intrinsic properties satisfying all these requirements. Coupling materials with synergic properties could be the most suitable way to design and build efficient photoelectrodes<sup>5</sup>. For example, it has been reportedly shown

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that coupling cocatalysts  $^{6,7}$ , charge collectors  $^8$ , and homo- or heterojunctions  $^{9-12}$  can improve the photocatalytic performance of nanomaterials.

In the more recent years, the n-type  $\rm BiVO_4$  semiconductor has been widely studied as a photoanode in water-splitting PEC cells taking into account that it is composed of chemical elements that are relatively abundant on Earth, has a short bandgap energy (2.5 eV), its valence band is suitably positioned towards the water oxidation, and its conduction band is also suitably placed close to the  $\rm H_2$  evolution potential  $^{13}$ . However, the primary challenge for improving such photoactive  $\rm BiVO_4$  consists in minimizing the fast electron-hole recombination due to the slow electron transport in  $\rm BiVO_4$ . Consequently, extremely low photovoltages (typically,  $\rm V_{ph} < 0.6 \, V$ ) have been reported for  $\rm BiVO_4$   $^{14,15}$ .

The more common strategy to increase the photoactivity of  $BiVO_4$  is based on the formation of p-n heterojunctions. Various p-n heterojunctions such as  $BiOI/BiVO_4^{16,17}$ ,  $BiOCI/BiVO_4^{18}$ ,  $Bi_2O_3/BiVO_4^{19,20}$ ,  $NiO/BiVO_4^{21}$ ,  $Si/BiVO_4^{22}$ ,  $CuO/BiVO_4^{23}$ ,  $Cu_2O/BiVO_4^{24}$ ,  $Ag_2O/BiVO_4^{25}$ , and  $Co_3O_4/BiVO_4^{26}$  have been proven to be an efficient method for keeping electron-hole pairs separated due to the built-in potential produced by the heterojunction. However, those mentioned p-type semiconductors are either unstable in water under oxidizing conditions or unable to absorb a wider visible light range than  $BiVO_4$ , thus limiting its application on photoelectrochemical water splitting.

The water splitting process in PEC cells is driven by the band edge pinning effect, in which the band edge positions of a semiconductor are pinned about the water redox potential. Thus, the height of band bending and consequently, the  $V_{ph}$  can be estimated by the difference between the flatband potential ( $V_{fb}$ ) of the semiconductor and the chemical potential of the targeted reaction in water<sup>27</sup>. Regarding the BiVO<sub>4</sub>, the maximum theoretical photovoltage for the water oxidation is approximately 1.23 V based on its flatband potential of about 0  $V_{RHE}^{7,13}$ . However, that theoretical  $V_{ph}$  value has not been reached up to now. We understand that the surface electronic states result in the pinning of the Fermi level of BiVO<sub>4</sub>, which decreases the experimental  $V_{ph}$  about the theoretical value<sup>28</sup>. To solve this, we can introduce internal electrical fields through ferroelectric polarization method by forming buried semiconductor/ferroelectric semiconductor heterojunctions.

For this purpose, in this study we made a heterojunction by coupling  $BiVO_4$  with  $Bi_4V_2O_{11}$  (Fig. 1), which is an n-type layered semiconductor that exhibits strong polar responses due to its polarized oxygen deficient perovskite slabs (i.e.,  $(VO_{3.5\,0.5})^{2-}$ ; = anionic vacancy)<sup>29,30</sup>. When the ferroelectric  $Bi_4V_2O_{11}$  is coupled to the  $BiVO_4$ , its spontaneous electrical polarization generates a sufficiently high electric field that produces an inversion layer at the  $BiVO_4/Bi_4V_2O_{11}$  interface (Fig. 1a), thus resulting in an increased band bending, enhancing the charge generation and the separation efficiency. By this approach,  $BiVO_4$  is the light-absorbing material for generating the electron-hole pairs, whereas the ferroelectric  $Bi_4V_2O_{11}$  perovskite plays the role of absorbing light and creating a built-in electric field to separate the photogenerated charges, inducing the carrier transport. A key advantage of the  $BiVO_4/Bi_4V_2O_{11}$  coupling junction is the appropriate energy band alignment (Fig. 1b), which facilitates the current flow<sup>31</sup>. We have shown that the photovoltage was readily tuned by doping the  $BiVO_4/Bi_4V_2O_{11}$  heterojunction with tungsten ions. As a result, the photovoltage and photocurrent were enhanced as much as 83% and 180%, respectively.

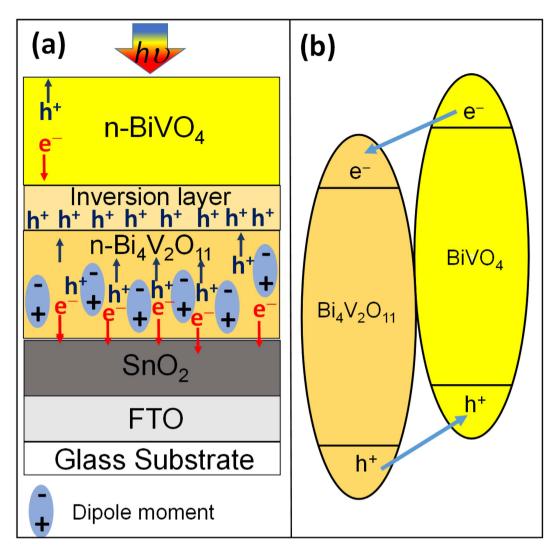
## Results

**Photoelectrode synthesis and characterization.** The XRD pattern of the sample  $Bi_4V_2O_{11}$  (Fig. S1a) confirmed the presence of  $Bi_4V_2O_{11}$  (a=5.5479(5) Å, b=5.5522(1) Å, and c=15.4665(2) Å) without any impurity, while characteristics peaks of  $BiVO_4$  (a=5.1808(2) Å, b=5.0915(1) Å and c=11.6801(1) Å) could be observed in Fig. S1b. The  $SnO_2$  peaks seen in the XRD patterns of these samples are due to fluorine-doped tin oxide (FTO). The  $BiVO_4/Bi_4V_2O_{11}$  heterojunction prepared through the Pechini method was deposited onto FTO glass covered with  $SnO_2$  to act as a passivating layer by using the spray pyrolysis method as previously elsewhere reported<sup>32</sup>. The XRD pattern of the sample labeled W0 (sample of the undoped heterojunction) (Fig. S1c) revealed that the as-prepared film is constituted by the monoclinic scheelite  $BiVO_4$  phase (a=5.1811(2) Å, b=5.0912(1) Å and c=11.6807(1) Å) and the orthorhombic  $Bi_4V_2O_{11}$  (a=5.5488(5) Å, b=5.5520(1) Å, and c=15.4670(2) Å). To increase the photovoltage and facilitate the photocurrent, the tungsten was used as a dopant into the  $BiVO_4/Bi_4V_2O_{11}$  film (sample labeled W1)<sup>33</sup>. The successful structural doping with W into the films was confirmed by the Rietveld analysis of the XRD data, which indicated the expanded unit cells for the W-doped  $BiVO_4/Bi_4V_2O_{11}$  (a=5.84873(9) Å) (Fig. S1d). The quantitative Rietveld analysis showed that both heterojunctions are formed by 52 wt.%  $BiVO_4$  and 48 wt.%  $Bi_4V_2O_{11}$ .

SEM images of the heterojunctions (Fig. S2) revealed the formation of  $BiVO_4/Bi_4V_2O_{11}$  islands with a highly porous fiber-like morphology. The thickness of both heterojunction films determined by cross-sectional SEM was  $9\,\mu m$  (Fig. S3). The EDS mapping (Fig. S4) showed the uniformly distributed Bi and V elements throughout the films. In addition to Bi and V, it was also detected a homogeneous distribution of W all over the W1 heterojunction. The W content in the film was in average 1.2(3) wt.%, as determined by XRF analysis.

The DRS spectra (Fig. S5a) showed that the films absorb radiation over a wide range of wavelengths, from the UV to the visible regions of the optical spectrum. The optical band gap energies (Fig. S5b) estimated for the  $BiVO_4$  and  $Bi_4V_2O_{11}$  in the undoped heterojunction were 2.53 and 2.42 eV, respectively. This means that both components of the heterojunction can be excited by visible light to produce reactive species for the water splitting. The band gap energies of  $BiVO_4$  and  $Bi_4V_2O_{11}$  in the W-doped sample (W1 in the Fig. S5) are similar to those of the non-doped heterojunction. No visible impurity absorption exists as the W doping energy levels are shallow<sup>34</sup>.

**Photoelectrochemical performance.** The current-potential curves presented in Fig. 2 clearly shows that the W doping of the BiVO<sub>4</sub>/Bi<sub>4</sub>V<sub>2</sub>O<sub>11</sub> heterojunction improved the observed photocurrent of the photoelectrode under visible light ( $\lambda > 450 \text{ nm}$ ). Photocurrent values at 1.23 V<sub>RHE</sub> were 0.05 and 0.14 mA cm<sup>-2</sup>



**Figure 1.** (a) The internal electrical field generated by the ferroelectric  $Bi_4V_2O_{11}$ . Schematic drawing of n-BiVO<sub>4</sub>/ferroelectric n-Bi<sub>4</sub>V<sub>2</sub>O<sub>11</sub> heterojunction with a polarization-induced inversion layer, and (b) Band alignment. Schematic illustration of BiVO<sub>4</sub>/Bi<sub>4</sub>V<sub>2</sub>O<sub>11</sub> band alignment.

for the W0 and W1 heterojunctions, respectively. It has been reported that the replacement of V sites by W as dopant eliminates the hole traps and extend the electron lifetime, thus significantly increasing the photocurrent<sup>33</sup>. Indeed, the longer electron lifetime reflects the slower recombination process in W1 than W0 film (Fig. S6). Interestingly, the photocurrent values obtained with the BiVO<sub>4</sub>/Bi<sub>4</sub>V<sub>2</sub>O<sub>11</sub> heterojunctions were significantly higher than those with pristine BiVO $_4$  (0.012 mA cm $^{-2}$ ) or Bi $_4$ V $_2$ O $_{11}$  (9  $\times$  10 $^{-5}$  mA cm $^{-2}$ ) at 1.23 V vs. RHE (Fig. 2), indicating that the heterojunction formation enhances the charge separation and consequently the photocurrent. Corroborating to these data, the electrochemical impedance analysis (Fig. S7) showed that the hole-transfer resistance across the heterojunction-electrolyte interface is significantly smaller than those of pristine BiVO<sub>4</sub> and Bi<sub>4</sub>V<sub>2</sub>O<sub>11</sub>, pointing out that the charge transfer processes on the BiVO<sub>4</sub>/Bi<sub>4</sub>V<sub>2</sub>O<sub>11</sub> heterojunction are faster than those occurring on the surface of the individual components. Additionally, the charge transfer resistance across the W-doped heterojunction-electrolyte interface is approximately one-seventh relatively to the non-doped heterojunction, implying that the W doping can also significantly improve the charge transfer across the heterojunction-electrolyte interface. In other words, the charge transfer from the near surface region of the W-doped BiVO<sub>4</sub>/Bi<sub>4</sub>V<sub>2</sub>O<sub>11</sub> heterojunction to the redox species on the surface is kinetically faster. Thus, as the charge transfer process becomes faster, it can efficiently minimize the electron-hole recombination at the BiVO<sub>4</sub>/ Bi<sub>4</sub>V<sub>2</sub>O<sub>11</sub> surface. As a result, the half-cell solar-to-hydrogen efficiency (HC-STH)<sup>35</sup> in a three-electrode setup was substantially increased after the heterojunction formation and W-doping treatment. The HC-STH values at  $0.85\,V_{RHE}$  for the bare BiVO<sub>4</sub>, Bi<sub>4</sub>V<sub>2</sub>O<sub>11</sub>, W0, and W1 films were 0.06, 0.00, 0.27 and 0.60%, respectively (Fig. S8). The incident photon-to-current efficiency (IPCE) spectra in a two-electrode setup for front-side illuminated W0 and W1 heterojunction at the potential of 0 V<sub>Pt</sub> (Fig. 3) showed a steady increase from 500 to 350 nm. The IPCE for the W1 film reached 2.6% at 450 nm with a maximum of 4.5% at 350 nm. On the other hand, the IPCE for the W0 film was significantly lower, reaching 1.0% at 450 nm and 1.4% at 340 nm at  $0 V_{Pt}$ .

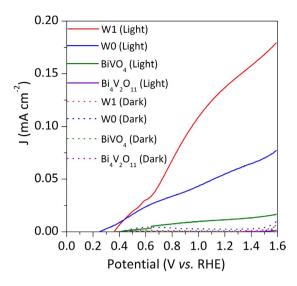
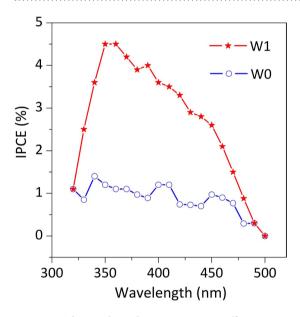


Figure 2. Photoelectrochemical characteristic of the photoelectrodes. Current-potential curves for the W0 and W1 films. Measurement conditions: active area of  $1.0\,\mathrm{cm^2}$ ,  $0.5\,\mathrm{M}$  Na<sub>2</sub>SO<sub>4</sub> electrolyte (pH = 6.6), Light Source: White Light ( $\lambda > 450\,\mathrm{nm}$ ,  $5\,\mathrm{mW}$  cm<sup>-2</sup>), scan rate of  $20\,\mathrm{mV}$  s<sup>-1</sup> from low to high potential, front illumination for both photoanodes.



**Figure 3.** The incident photon-to-current efficiency. IPCE spectra of the W0 and W1 heterojunctions at 0 V vs. Pt in  $0.5\,\mathrm{M}\,\mathrm{Na_2SO_4}$  electrolyte.

**Open-circuit potential measurements.** To better understand the effect of the heterojunction and W as a doping element on the enhancement of the photocurrent, we have tried to estimate the onset potential  $(V_{on})$  for each photoelectrode<sup>36</sup>. The  $V_{on}$  value for the pristine BiVO<sub>4</sub> was  $0.50\,V_{RHE}$ , whereas the heterojunction formation and W doping led to an increase of the  $V_{on}$  value. The  $V_{on}$  of  $0.54\,V_{RHE}$  was obtained for the W0 photoelectrode; after the W-doping, the  $V_{on}$  reached  $0.66\,V_{RHE}$  (Fig. S9). These trends are opposite to those obtained by electrochemical impedance (Fig. S7), which clearly showed that the W-doping into BiVO<sub>4</sub>/Bi<sub>4</sub>V<sub>2</sub>O<sub>11</sub> significantly enhances the kinetics of charge transfer across the semiconductor-electrolyte interface by reducing the kinetic overpotentials. It suggests that the current vs. potential measurements are not suitable for quantitative analysis, as the change in  $V_{on}$  may be affected for both kinetics (reduction in kinetics overpotentials) and thermodynamic (increase in  $V_{ph}$ ) factors<sup>37</sup>. To evaluate which factor plays a more important role on improving the photocurrent, we then measured the open-circuit potential (OCP) in the dark and under visible light. It enabled us to probe the resting potentials at zero net exchange photocurrents, where the kinetic factors are negligible<sup>37</sup>. The difference between the measured potential in light and that in the dark led to the  $V_{ph}$ . Hence, the pristine BiVO<sub>4</sub> and Bi<sub>4</sub>V<sub>2</sub>O<sub>11</sub> produced photovoltages of 0.11 and 0.13 V respectively, whereas the W0 and W1 heterojunctions gave the significantly higher photovoltages of 0.76 and 1.39 V, respectively (Fig. 4). The enlarged band bending at

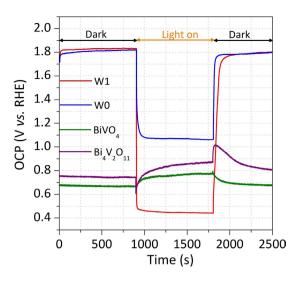


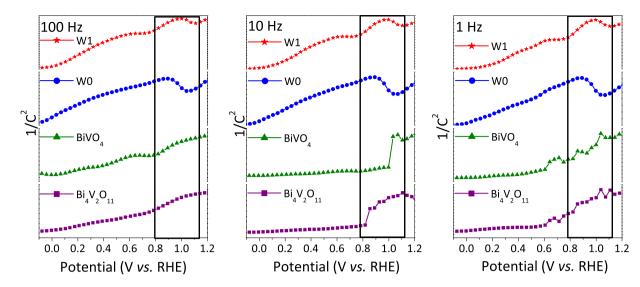
Figure 4. Photovoltage generated by the photoelectrodes. Open circuit potential measurements of BiVO<sub>4</sub>, Bi<sub>4</sub>V<sub>2</sub>O<sub>11</sub>, W0, and W1 photoelectrodes. Measurement conditions: active area of  $1.1\,\mathrm{cm}^2$ ,  $0.5\,\mathrm{M}$  Na<sub>2</sub>SO<sub>4</sub> electrolyte (pH = 6.6), light source: white light ( $\lambda$  > 450 nm, 5 mW cm<sup>-2</sup>).

the  $BiVO_4/Bi_4V_2O_{11}$  heterojunctions represent the enhanced electron-hole separation in the photoelectrodes. These data suggest that the heterojunction formation is essential to obtain high photovoltages. Moreover, the significant increase of 83% of the photovoltage value, compared to the non-doped heterojunction suggests that the W-doping treatment changes the  $BiVO_4/Bi_4V_2O_{11}$  interface, thus increasing the degree of band bending. It is worth to note that the OCP in the dark usually reflects the upward band bending nature of the photoanode in dark equilibrium with the electrolyte. However, the significantly more positive  $OCP_{dark}$  values (1.8 V vs. RHE) were obtained with the heterojunctions, suggesting that the surface states of the  $BiVO_4/Bi_4V_2O_{11}$  heterojunction were passivated to unpin the Fermi levels. As a result, an enlarged band bending was produced after the heterojunction formation. Furthermore, the obtained photovoltages were much higher than that predicted by the Schottky-Mott rule (0.61 V)<sup>38</sup> owing to the formation of an interface dipole at the n-n  $BiVO_4/Bi_4V_2O_{11}$  heterojunction, suggesting that the photovoltage generation by a simple n-n heterojunction is not the primary process governing the photovoltage.

**Mott-Schottky data.** To gain more insights into the factors that improved the photovoltage, we investigated the electronic characteristics of the BiVO<sub>4</sub>/Bi<sub>4</sub>V<sub>2</sub>O<sub>11</sub> heterojunction using electrochemical impedance spectroscopic, and the Mott-Schottky equation was used to evaluate the type of conductivity of the heterojunctions. Figure 5 shows the capacitance values of the space charge region obtained at various frequencies (1, 10 and 100 Hz). According to the Mott-Schottky equation, a linear relationship of 1/C<sup>2</sup> versus applied potential achieved in Fig. 5 displayed straight lines with positive slopes for all the studied frequencies, which corresponds to depletion regions typical of n-type semiconductors. In this case, the photogenerated holes can be spontaneously driven to the BiVO<sub>4</sub>/Bi<sub>4</sub>V<sub>2</sub>O<sub>11</sub> surface by the band bending in the space-charge layer, that is, where the electrode can act as a photoanode. Remarkably, in addition to the positive slopes, a region with a negative slope at sufficiently anodic voltages of the flatband potential was also observed in the W0 and W1 heterojunctions. It corresponds to a positively charged inversion layer, at which holes are the majority at the BiVO<sub>4</sub>/Bi<sub>4</sub>V<sub>2</sub>O<sub>11</sub> interface. The inversion layer phenomenon was not detected in the single BiVO<sub>4</sub> and Bi<sub>4</sub>V<sub>2</sub>O<sub>11</sub>, which suggests that this process may be occurring at the heterojunction interface. It reveals that in the inverted layer, the BiVO<sub>4</sub>/Bi<sub>4</sub>V<sub>2</sub>O<sub>11</sub> heterojunctions can work as a photocathode. Therefore, the behavior of the films as photocathodes was also analyzed. Indeed, it was verified through the current-potential curves (Fig. S10) that both non-doped and W-doped BiVO<sub>4</sub>/Bi<sub>4</sub>V<sub>2</sub>O<sub>11</sub> heterojunctions exhibited photoresponses that are typical of p-type semiconductors at cathodic potentials. The photocurrent values at  $0\,V_{RHE}$  for the W0 and W1 films were -0.03 and  $-0.14\,\text{mA}$  cm $^{-2}$ , respectively. On the other hand, the pristine Bi<sub>4</sub>V<sub>2</sub>O<sub>11</sub> and BiVO<sub>4</sub> exhibited smaller photocurrent values of -0.004 and -0.010 mA cm<sup>-2</sup> respectively, indicating that the p-type behavior of the photoelectrodes is typical of the heterojunction.

### Discussion

Given the above results, we suggest that the high photovoltages generated by the  $BiVO_4/Bi_4V_2O_{11}$  heterojunctions may be assigned to the strong inversion layer at the  $BiVO_4/Bi_4V_2O_{11}$  interface. The self-polarization due to the existence of  $(VO_{3.5\,0.5})^{2-}$  layers in  $Bi_4V_2O_{11}$  creates a sufficiently high internal polar electric field to produce a hole inversion layer at the  $BiVO_4/Bi_4V_2O_{11}$  interface, as verified from the Mott-Schottky data. The internal surface dipole on the  $Bi_4V_2O_{11}$  increases the local vacuum level at the heterojunction interface, thus changing the energy needed to extract an electron from the semiconductor into the vacuum and, therefore, can be thought of as altering the ionization potential, electron affinity, and the work function of the semiconductor. When the  $Bi_4V_2O_{11}$  containing dipole moments is brought into intimate contact with the  $BiVO_4$ , it is the difference between the effective work functions of the two that is the driving force for achieving electronic equilibrium via charge transfer between them and the creation of a space charge region within the heterojunction. Thus, charges of the same



**Figure 5. Inversion layer detected by Mott-Schottky plots.** Variation of capacitance (C) with the applied potential in  $0.5\,\mathrm{M}\,\mathrm{Na_2SO_4}$  (pH = 6.6) presented in the Mott-Schottky relationship for the films BiVO<sub>4</sub>, Bi<sub>4</sub>V<sub>2</sub>O<sub>11</sub>, W0, and W1 at different frequencies.

polarity will be repelled from the heterojunction interface whereas those of the opposite polarity will be attracted to it (Fig. 1). Hence, if the ferroelectric  $Bi_4V_2O_{11}$  is in contact with the n-type  $BiVO_4$ , majority electrons will be depleted from the  $BiVO_4/Bi_4V_2O_{11}$  interface, while minority holes will be attracted to it, thus inverting its conductivity to  $p^+$  ( $p^+$  denotes highly doped semiconductors compared with p used for low/moderately doped one). As a consequence, a virtual p-n junction similar to the p-n junction occurring in a metal-insulating-semiconductor (MIS)<sup>39</sup>, can be observed. This is in agreement with the photoanodic and photocathodic behavior exhibited by the  $BiVO_4/Bi_4V_2O_{11}$  heterojunction and Mott-Schottky data, which revealed typical characteristics of n- and p-type semiconductors. Because of that, photogenerated charges can be more efficiently separated by the virtual p-n heterojunction, thus providing high photovoltages. It is worth to note that the net photovoltage strongly depends on the polarization in  $Bi_4V_2O_{11}$ , which can be efficiently tuned by creating oxygen vacancies and strain through W-doping into its structure. Additionally, the inversion layer has been shown to decrease the impact of surface trap states, thus reducing the surface recombination and consequently increasing the observed photocurrent<sup>40</sup>.

In summary, this work addresses a central topic on water splitting into PEC cells related to the low photovoltages generated by semiconductors. We have shown that the self-polarization of  $Bi_4V_2O_{11}$  perovskite creates an internal electrical field, which induces the formation of a hole inversion layer at the  $BiVO_4/Bi_4V_2O_{11}$  interface. Because of that, a virtual p-n junction was formed at the  $BiVO_4/Bi_4V_2O_{11}$  interface and consequently a high photovoltage was produced. The self-polarization effect was tuned by doping the ferroelectric material with W and, therefore, an increase of 83% in the generated photovoltage was developed as compared to the undoped heterojunction. Because the high photovoltage can promote an enhanced charge separation at the photoelectrodes, an increase of 180% in the photocurrent compared to the undoped heterojunction was achieved. We hope that the approach described here may enable the design of practical photoelectrodes based on earth-abundant semiconductors with high photovoltage and enhanced charge separation for water splitting.

#### Methods

**Preparation of the precursor solution.** All chemicals were used as received without further purification. Ten millimoles of ammonium metavanadate ( $NH_4VO_3$ , 99%) were dissolved in 50 mL of 174.2 mmol  $NH_4OH$  (24.5%) to get the solution "A". In another recipient, 10 mmol of  $Bi(NO_3)_3$ · $5H_2O$  was dissolved in 50 mL of 174.2 mmol  $CH_3COOH$  (99.7%), producing the solution "B". Both the solutions were stirred separately under heating at 80 °C to form stable and homogeneous solutions. Then, the solution A and B were mixed under stirring at 80 °C for 90 min. Next, 40 mmol citric acid was added to the mixture, and the resultant solution was kept standing for 24 h. The solution was then diluted with 100 mL deionized water and heated at 80 °C until the solution became light blue. This homogeneous solution was used as a precursor for preparing the photoelectrodes. Tungsten doped sample was achieved by a similar way as described above, except by the addition of ammonium metatungstate solution to the solution "A".

**Preparation of the photoelectrodes.** Transparent conductive FTO-coated glass ( $10\,\mathrm{mm}\times30\,\mathrm{mm}\times2\,\mathrm{mm}$ ,  $16\,\Omega\,\mathrm{cm}^{-2}$ ) were used as substrates for the deposition of the films. Before the films deposition onto FTO, the substrate surface was cleaned in an ultrasonic bath using acetone and ethanol each for  $15\,\mathrm{min}$ . After the cleaning procedure, the substrates were dried in a furnace at  $120\,^{\circ}\mathrm{C}$  for  $1\,\mathrm{h}$ . Then, a  $\mathrm{SnO}_2$  passivating layer was deposited onto FTO. Briefly,  $0.4513\,\mathrm{g}\,\mathrm{SnCl}_2\cdot\mathrm{2H}_2\mathrm{O}$  was dissolved in  $50\,\mathrm{mL}$  ethanol (99.8%) and 2 drops of  $\mathrm{CH}_3\mathrm{COOH}$  (99.7%). This solution was sprayed down using a commercial airbrush ( $0.3\,\mathrm{mm}$  nozzle) directly onto the FTO glass onto FTO at  $300\,^{\circ}\mathrm{C}$  at a distance of  $20\,\mathrm{cm}$ , with 2 cycles of spray deposition for  $5\,\mathrm{s}$  each cycle. Films were subsequently annealed in air in a muffle furnace at  $450\,^{\circ}\mathrm{C}$  for  $2\,\mathrm{h}$ , to produce a  $\mathrm{SnO}_2$  layer. Finally, the precursor

solutions were sprayed down onto the FTO glass at a temperature of 300 °C at a distance of 20 cm. 1 cycle of spray with 5 s time deposition was used. The films were subsequently annealed in air in a muffle furnace at 500 °C for 5 h to produce the undoped  $BiVO_4/Bi_4V_2O_{11}$  (sample W0) and W-doped  $BiVO_4/Bi_4V_2O_{11}$  (sample W1) photoelectrodes.

**Characterization.** The morphology of films was investigated by scanning electron microscopy (SEM) using a tabletop SEM (Hitachi TM - 300). EDS mapping was obtained in a SwiftED3000 (Oxford Instruments) at 15 kV accelerating voltage. The crystalline phases of the films were determined using an X-ray diffractometer (XRD 6000, Shimadzu). The data were collected from 10 to 80° 20 at a step width of 0.2°, 10 s per step, at 40 kV, 200 mA, and CuK $\alpha$  radiation ( $\lambda$  = 1.540560 Å). Silicon was used as an external standard. The Rietveld structural refinement was performed with FullProf\_Suite 2015 software. The diffuse reflectance spectra were collected with a UV-Vis spectrophotometer (Shimadzu UV 2700). Teflon was used as reference material (100% transmission). The direct bandgap energies were calculated by the following Tauc equation (1):

$$(\alpha h \nu)^2 = A(h \nu - E_g) \tag{1}$$

where A = constant,  $h\nu$  = light energy,  $E_g$  = optical bandgap energy,  $\alpha$  = measured absorption coefficient.

**Photoelectrochemical measurements.** The photoelectrochemical measurements were carried out with a potentiostat (AUTOLAB Potentiostat-Galvanostat PGSTAT 128 N) using a standard three-electrode cell with an Ag/AgCl (3.0 M KCl) reference electrode, a platinum wire as a counter electrode, a working electrode with irradiation area of  $1.1~\rm cm^2$ , and a scan rate of  $20~\rm mV~s^{-1}$ . A  $0.5~\rm M~Na_2SO_4$  aqueous solution (pH = 6.6) was used as electrolyte. The prepared films were connected to a copper tape to measure the photoactivity. The current-potential curves were recorded in the dark and under a white light LED (light intensity of  $5~\rm mW~cm^{-2}$ ,  $\lambda > 450~\rm nm$ ). For converting the obtained potential vs. Ag/AgCl to RHE, the equation (2) was used:

$$E_{RHE} = E_{Ag/AgCl} + 0.059pH + E_{Ag/AgCl}^{o}$$
 (2)

The  $E^{o}_{Ag/AgCl}$  (3 M KCl) = 0.197 at 25 °C.

**Electrochemical impedance spectroscopy.** The electrochemical impedance spectroscopy was performed using an AUTOLAB Potentiostat-Galvanostat PGSTAT 128 N equipped with the FRA32M module. The Nyquist plots were measured at 0.7 V vs. Ag/AgCl with an AC amplitude of 20 mV, frequency of 100 kHz—100 mHz under a white light LED ((light intensity of 5 mW cm<sup>-2</sup>,  $\lambda$  > 450 nm). The measured spectra were fitted using the NOVA 1.11 software. The 0.5 M Na<sub>2</sub>SO<sub>4</sub> solution was utilized for the all electrochemical measurements.

Mott-Schottky data acquisition is a usual way for the electrochemical characterization of semiconductor materials. The spectra were collected using an AUTOLAB Potentiostat-Galvanostat PGSTAT 128 N equipped with the FRA32M module, and a cell configuration with three electrodes: a reference electrode of Ag/AgCl (3, 0 M KCl), a platinum as counter electrode wire and a working electrode (W0 or W1) with  $1.0~\rm cm^2$  irradiation area under white light LED (light intensity of 5 mW cm<sup>-2</sup>,  $\lambda$  > 450 nm), applying potential  $-0.6~\rm V$  to  $+0.6~\rm V$  vs. Ag/AgCl in the frequency range of 1 to 100 Hz. The measured spectra were fitted using the NOVA 1.11 software. The 0.5 M Na<sub>2</sub>SO<sub>4</sub> solution was utilized for the all electrochemical measurements.

**IPCE measurements.** Incident photon-to-current efficiency (IPCE) spectra were determined from photocurrent measurements, and the equation (3) was used:

$$IPCE = \frac{(1240 \text{ eV nm})I_{\text{ph}}}{\lambda P_{\text{irr}}}$$
(3)

where,  $I_{ph}$  is the photocurrent density in mA cm $^{-2}$ , l is the irradiation wavelength of the incident light in nm, and  $P_{irr}$  is the incident photon flux in mW cm $^{-2}$ . In these experiments, the excitation was obtained by a 300 W xenon lamp in a Newport/Oriel housing, which was passed through a Newport/Oriel Cornerstone 260 monochromator to produce monochromatic light with a bandwidth of 10 nm. Photon flux was measured with a Newport 1919-R power meter equipped with an 818-UV/DB sensor before and after photocurrent measurements to ensure stability and reproducibility. Photo-action spectra were acquired at 10 nm intervals between 320 and 500 nm for the bared and W-doped BiVO $_4/Bi_4V_2O_{11}$  films immersed in 0.5 M  $Na_2SO_4$  solution and using a platinized FTO as the counter electrode. The photoanode and the counter-electrode were mounted in a sandwich-like setup, using an inert polymer (50  $\mu$ m Dupont Surlyn) as a spacer. The space between the electrodes was fulfilled by the electrolyte using a vacuum.

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# **Author Contributions**

M.C.P. and J.D.F. supervised the project and wrote the manuscript. W.S.S. and M.R. planned, performed and interpreted the photoelectrochemical measurements. A.S.A. and A.C.S. performed and interpreted the electrochemical impedance data. J.P.M. carried out the XRD and SEM measurements. L.L.N. collected the IPCE data. A.O.T.P. interpreted the IPCE data. L.C.A.O. performed and discussed the DRS data. All authors discussed the results and commented on the manuscript.

# **Additional Information**

**Supplementary information** accompanies this paper at http://www.nature.com/srep

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